



Swansea University
Prifysgol Abertawe



Cronfa - Swansea University Open Access Repository

This is an author produced version of a paper published in :

Nano Letters

Cronfa URL for this paper:

<http://cronfa.swan.ac.uk/Record/cronfa32365>

Paper:

Carrad, D., Mostert, A., Ullah, A., Burke, A., Joyce, H., Tan, H., Jagadish, C., Krogstrup, P., Nygård, J., Meredith, P. & Micolich, A. (2016). Hybrid Nanowire Ion-to-Electron Transducers for Integrated Bioelectronic Circuitry. *Nano Letters*, 17(2), 827-833.

<http://dx.doi.org/10.1021/acs.nanolett.6b04075>

This article is brought to you by Swansea University. Any person downloading material is agreeing to abide by the terms of the repository licence. Authors are personally responsible for adhering to publisher restrictions or conditions. When uploading content they are required to comply with their publisher agreement and the SHERPA RoMEO database to judge whether or not it is copyright safe to add this version of the paper to this repository.

<http://www.swansea.ac.uk/iss/researchsupport/cronfa-support/>

Hybrid nanowire ion-to-electron transducers for integrated bioelectronic circuitry

D.J. Carrad,[†] A.B. Mostert,[¶] A.R. Ullah,[†] A.M. Burke,[†] H.J. Joyce,^{||} H.H. Tan,[#]
C. Jagadish,[#] P. Krogstrup,[@] J. Nygård,[@] P. Meredith,[¶] and A.P. Micolich^{*,†}

School of Physics, University of New South Wales, Sydney NSW 2052, Australia, Walter Schottky Institut, Technische Universität München, Am Coulombwall 4, Garching 85748, Germany, Centre for Organic Photonics and Electronics, School of Mathematics and Physics, University of Queensland, Brisbane QLD 4072, Australia, Solid State Physics/NanoLund, Lund University, SE-221 00 Lund, Sweden, Department of Electronic Materials Engineering, Research School of Physics and Engineering, The Australian National University, Canberra ACT 2601, Australia, Department of Engineering, University of Cambridge, Cambridge CB3 0FA, United Kingdom, Department of Electronic Materials Engineering, Research School of Physics and Engineering, The Australian National University, Canberra ACT 0200, Australia, Center for Quantum Devices, Nano-Science Center, Niels Bohr Institute, University of Copenhagen, Universitetsparken 5, DK-2100 Copenhagen, Denmark, and Physics Department, Swansea University, Swansea SA2 8PP, Wales, United Kingdom

E-mail: adam.micolich@nanoelectronics.physics.unsw.edu.au

Abstract

A key task in the emerging field of bioelectronics is the transduction between ionic/protonic and electronic signals at high fidelity. This is a considerable challenge since the two carrier types exhibit intrinsically different physics and are best supported by very different materials types – electronic signals in inorganic semiconductors and ionic/protonic signals in organic or bio-organic polymers, gels or electrolytes. Here we demonstrate a new class of organic-inorganic transducing interface featuring semiconducting nanowires electrostatically gated using a solid proton-transporting hygroscopic polymer. This model platform allows us to study the basic transducing mechanisms as well as deliver high fidelity signal conversion by tapping into and drawing together the best candidates from traditionally disparate realms of electronic materials research. By combining complementary n - and p -type transducers we demonstrate functional logic with significant potential for scaling towards high-density integrated bioelectronic circuitry.

Keywords: III-V nanowires, bioelectronics, proton-to-electron transduction, hybrid organic/inorganic electronics.

Ion-to-electron signal transduction is at the heart of modern bioelectronics. To date, the predominant means to affect such transduction has been via organic semiconductor-based devices seeking to exploit the biocompatibility and processing flexibility of these materials.

*To whom correspondence should be addressed

†School of Physics, University of New South Wales, Sydney NSW 2052, Australia

‡Walter Schottky Institut, Technische Universität München, Am Coulombwall 4, Garching 85748, Germany

¶Centre for Organic Photonics and Electronics, School of Mathematics and Physics, University of Queensland, Brisbane QLD 4072, Australia

§Solid State Physics/NanoLund, Lund University, SE-221 00 Lund, Sweden

||Department of Electronic Materials Engineering, Research School of Physics and Engineering, The Australian National University, Canberra ACT 2601, Australia

⊥Department of Engineering, University of Cambridge, Cambridge CB3 0FA, United Kingdom

#Department of Electronic Materials Engineering, Research School of Physics and Engineering, The Australian National University, Canberra ACT 0200, Australia

©Center for Quantum Devices, Nano-Science Center, Niels Bohr Institute, University of Copenhagen, Universitetsparken 5, DK-2100 Copenhagen, Denmark

△Physics Department, Swansea University, Swansea SA2 8PP, Wales, United Kingdom

Typically these devices are transistors where changes in ionic environment affect the electric current through a semiconducting channel via an electrostatic field (organic field-effect transistors) or by ions entering the semiconductor and chemically doping or de-doping it (organic electrochemical transistors).^{1,2} The latter have shown particular promise for applications such as *in vivo* recording of brain activity.³ ‘Protonic’ transistors featuring thin proton conductive films bridging PdH_x electrodes^{4,5} have also been developed towards this application. A key long-term goal is full-scale integration of these devices with conventional electronics to facilitate signal recording, data transmission and control of multiplexed transducer arrays.⁶ An important first step is the development of ionically-active amplifiers and logic-gates featuring complementary *n*- and *p*-type transistors. Fully ionic junction transistors⁷ and logic gates⁸ have been developed, as have complementary protonic devices,⁹ but they face difficulties with slow switching, poorly defined off-states and/or low gain.

Here we report a new hybrid organic-inorganic materials platform with considerable scaling potential that presents an interesting route to overcoming these problems. We combine salt-free nanoscale polyethylene oxide (PEO) strips patterned by electron-beam lithography with *n*-type InAs and *p*-type GaAs nanowires to demonstrate proton-to-electron transduction at frequencies up to 50 Hz. This builds on previous work on polymer electrolyte-gated nanowire devices^{10,12} but improves their biocompatibility by removing the commonly added perchlorate salt. We also present functional prototype hybrid protonic/electronic complementary circuits, a first step towards direct biological signal amplification and bioelectronic logic.^{6,8} We attribute the transduction efficacy of our devices to the high surface-to-volume ratio of both the nanowire conducting channel and the patterned PEO dielectric, combined with the low area interface between them. In particular the high surface-to-volume ratio of our nanoscale PEO strips should drive significantly enhanced H₂O uptake compared to the thin-film parallel-plate capacitor geometries previously used to study ionic conductivity in PEO. This translated to a drastic improvement in PEO ionic conductivity and electrical performance. Strong ion-to-electron transduction can be obtained in inorganic semiconduc-

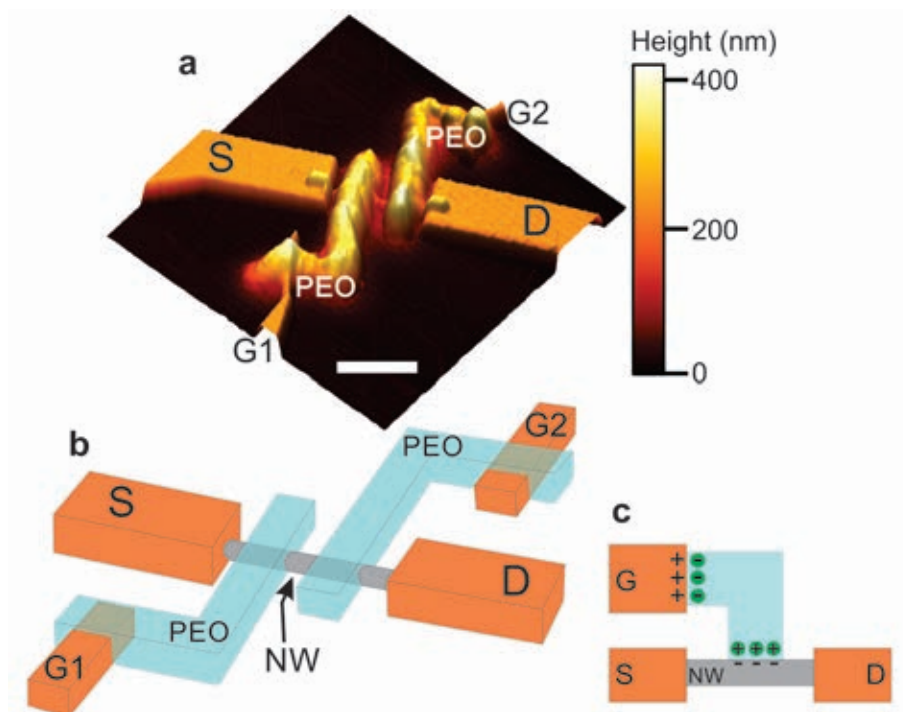


Figure 1: **PEO-gated nanowire field-effect transistor.** **a** Atomic force micrograph and **b** schematic showing the *n*-InAs or *p*-GaAs nanowire (NW) channel with metal source (S) and drain (D) contacts and gate electrodes (G1 and G2). The gates and nanowire were linked by electron-beam patterned PEO.¹⁰ The schematic in **c**) illustrates the separation of ionic species, e.g., H^+/OH^- , $\text{Li}^+/\text{ClO}_4^-$ across the PEO in response to an applied gate voltage. G1 and G2 were offset significantly from the channel here to ensure gating only occurs by proton drift and not by direct field-effect (see Supplementary Fig. S1).¹¹ The scale bar in **a** represents $2 \mu\text{m}$.

tor nanowires by virtue of their high surface-to-volume ratio and strongly surface-influenced conduction properties.^{13,14} These desirable properties have seen individual nanowires both with and without functionalized surfaces used to detect specific molecules^{13,14} and act as, e.g., gas¹⁵ and pH sensors.^{13,16} In our devices ionic separation across the PEO strip replaces polarization of the surface charge with the advantage that the ionic signal could potentially be simultaneously sensed by complementary transistors and thereby enable functional logic and amplification circuits. Additionally, by using an inorganic semiconductor channel we can tap into decades of materials optimisation and nanoscale integration knowledge enabling us to tailor our *n*- and *p*-type devices to the common operating voltage range necessary for functional complementary architectures. We chose III-V semiconductor nanowires over Si for their higher electron/hole mobility, noting also recent rapid improvements in direct integration of III-V nanowires with Si-based microelectronics.^{17,18}

Results

Device Design. Figure 1a/b shows the nanowire transistor structure used in these experiments. Two electron-beam patterned PEO strips^{10,19} extend from the end of metal gate electrodes (G1/G2) to cross an *n*-InAs or *p*-GaAs nanowire connected to source (S) and drain (D) contacts. Although in principle only one PEO strip is required to achieve a working device, our devices feature two for improved yield, to confirm proper strip isolation by independent gating,¹⁰ and to check reproducibility of performance. A gate voltage V_G on G1 or G2 drives ion drift within the PEO leading to electric double layers at the interfaces between the PEO and gate electrode and the PEO and nanowire as shown schematically in Fig. 1c. These electrical double layers effectively transfer the gate charge to within ~ 1 nm of the nanowire surface, as discussed by Kim *et al.*,²⁰ giving rise to a large gate capacitance and significantly reduced switching voltage. Typically the PEO is intentionally doped with a salt e.g. LiClO_4 , and $\text{Li}^+/\text{ClO}_4^-$ transport is assumed to dominate the gating effect. The

lack of salt-doping in our PEO and the device behaviour explored below instead point to H^+/OH^- as the active ionic species, consistent with the known protonic conductivity behaviour of undoped PEO.²¹ These H^+/OH^- ions are transported along PEO chains via the ethereal oxygen atoms exploiting local segmental motions of the polymer. We would also expect the Grotthuss mechanism of proton transfer along hydrogen-bonded networks in the adsorbed water to play a role, as reported recently for other proton-conductive polymers.^{4,9} Gating action at the nanowire surface will be a mixture of proton-accumulation field-effect and protonation/deprotonation of the nanowire’s native surface oxide, as occurs in InAs nanowire ion-sensitive field-effect transistors (ISFETs).^{13,14,16} The gate electrode terminates $\sim 4 \mu\text{m}$ from the nanowire, alongside the source or drain contact. This ‘offset’ gate geometry eliminates any competing influence by direct ‘vacuum’ field-effect or non-ionic dielectric polarisation.^{10,11} We do this to confirm ion transport is the dominant gating mechanism and to facilitate measurement of the PEO protonic conductivity; a detailed discussion is available in the Supplementary Information.

Gating action and humidity dependence of undoped polyethylene oxide. Figure 2 shows the gating action of our InAs and GaAs nanowire transistors with undoped PEO gate dielectrics. The InAs and GaAs devices in Fig. 2 exhibit clear *n*-type and *p*-type depletion, respectively. The devices are well matched, with subthreshold swings $\approx 300 \text{ mV/dec}$, on/off ratios $\approx 10^3 - 10^4$ and well overlapping operating voltage ranges $-1.5 < V_G < +1 \text{ V}$, thereby showing significant potential for complementary circuit architectures. The lower current I_d for the *p*-GaAs device on-state reflects the difficulty in making low resistance *p*-type ohmic contacts to GaAs nanowires; work continues towards improving this. We present data at three relative water vapour pressures $P_{rel} = 0, 0.5$ and 1 in Fig. 2; the full data set in $P_{rel} = 0.125$ increments is in the Supplementary Information. Hydration control was achieved using a specially constructed vacuum chamber connected to a reservoir of degassed, deionized H_2O .^{22,23} The existence of a hydration dependence is consistent with H_2O uptake, dissociation and H^+/OH^- transport being the key component of the gating action in our

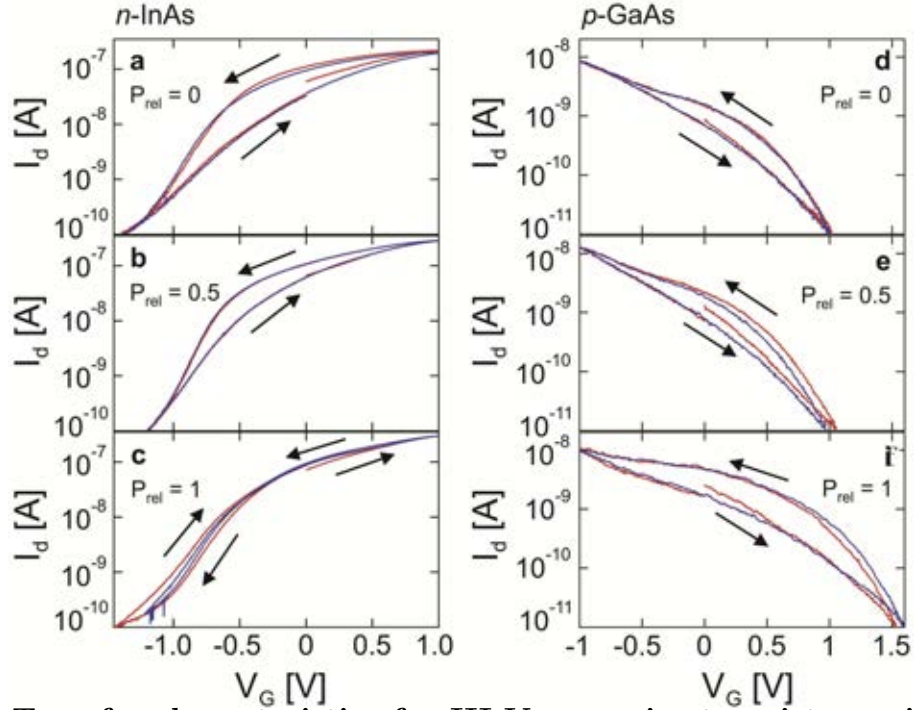


Figure 2: **Transfer characteristics for III-V nanowire transistors with PEO gate dielectrics.** Plots of nanowire channel current I_d vs polymer electrolyte gate voltage V_G for **a-c** an InAs nanowire transistor and **d-f** a GaAs nanowire transistor, both with undoped PEO gate dielectric. In each case we present two consecutive gate traces (red/blue) obtained at three different relative water vapour pressures P_{rel} . Trace direction is indicated by arrows, with sweeps towards depletion terminated just prior to reaching the full ‘off’ state to avoid the complex figure-of-eight pattern that emerges due to gate hysteresis.¹⁰ The lessening hysteresis with increased hydration is indicative of a proton-based gating mechanism. We present the entire data set in Supplementary Figs S2 and S3, which allows us to comment further on the gate hysteresis and the non-monotonic behaviour for p -GaAs devices.

devices. This behaviour is expected given the known high water uptake¹⁹ and evidence of proton conduction^{21,24} in PEO thin films (see Discussion for more details).

The most strongly affected aspect of our devices under humidity increase was device operating speed. This manifested experimentally as hysteresis in the quasi-dc gate characteristics and frequency limiting under ac switching operations. Given the importance of operating speed for potential signal transduction applications, we now discuss the two key contributions of charge trapping by nanowire surface-states and the undoped PEO’s proton conductivity. The gate hysteresis under quasi-dc conditions is substantially reduced at high P_{rel} for the n -InAs devices (Fig. 2a-c). This indicates the importance of H^+/OH^- transport in our devices and its increasing effect as H_2O adsorption and dissociation improves the proton conductivity. Some hysteresis remains at $P_{rel} = 1$, characteristic of InAs surface-state trapping.²⁵ The p -GaAs devices (Fig. 2g-i) show more complex trends with increasing P_{rel} . Hysteresis initially decreases – as for n -InAs devices – but increases again for higher $P_{rel} > 0.75$ (the trends are more apparent in the full data set shown in Supplementary Fig. S3). This likely indicates a stronger surface-state contribution for p -GaAs, consistent with literature for planar p -GaAs transistors.^{26,27} While the hysteresis is a feature of the quasi-dc characteristics, it has a lesser influence on the ac characteristics providing device operation is sufficiently fast that surface state trap occupancy remains quasi-static. One example is the GHz operation of InAs nanowire transistors²⁸ even when the quasi-dc characteristics exhibit hysteresis.²⁹ We now turn to data showing that surface states do not impede high fidelity transduction of biological/ionic signals.

ac signal response. The transistor response to ac gate signals that mimic biological signals is the most interesting aspect from a proton-to-electron transduction perspective. Figures 3a-h show the I_d response (black/left axis) to a square wave V_G input (red/right axis) for increasing frequency f at $P_{rel} = 1$. Each device is operated where I_d is linear in V_G , giving good square wave fidelity at $f = 1$ Hz (Fig. 3a/e). The slight I_d noise/instability in the on-state arises from surface-state effects during the square-wave maxima/minima.

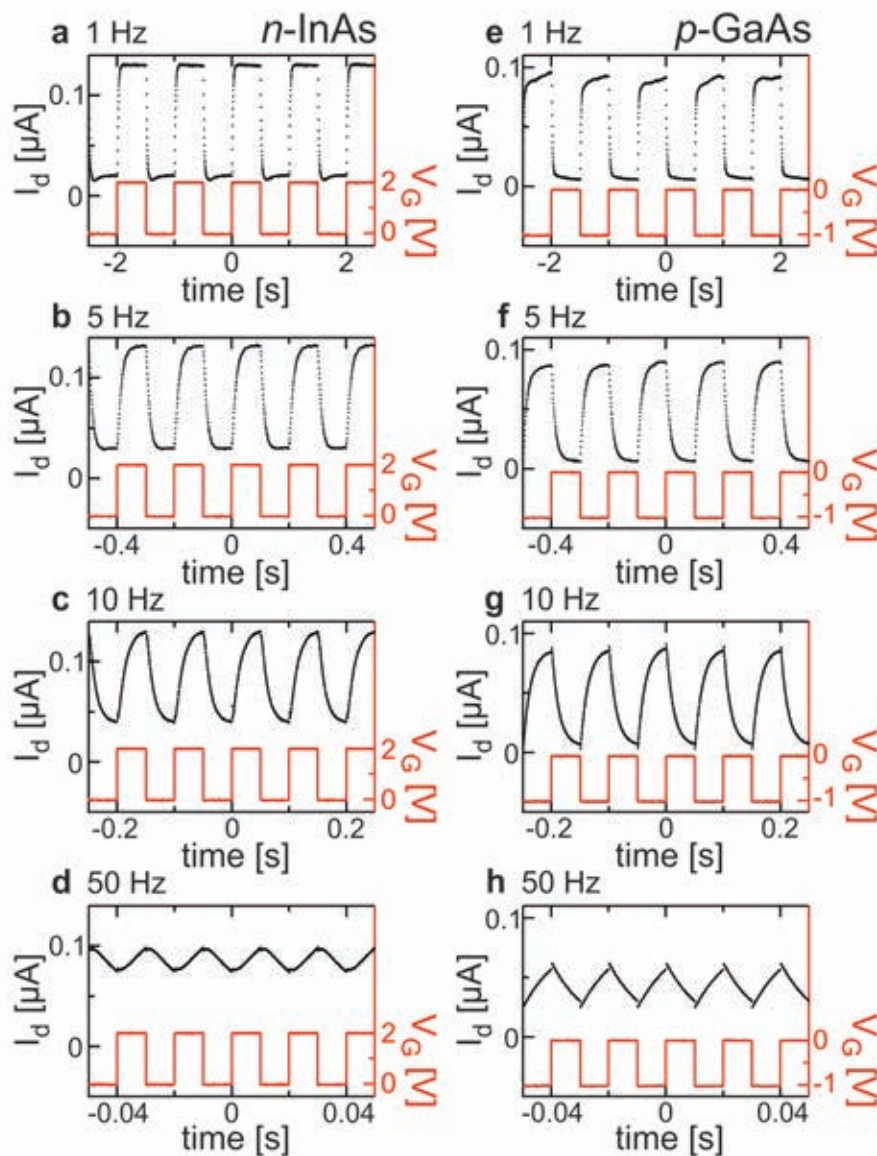


Figure 3: **Switching performance of NWFETs with PEO gate dielectrics at $P_{rel} = 1$.** **a-h** Plots of nanowire channel current I_d (black, left axis) vs time for a square wave applied to V_G (red, right axis) with frequency $f = 1, 5, 10$ and 50 Hz for **a-d** InAs and **e-h** GaAs nanowire channels. For this demonstration both devices were operated in a gate range where I_d is linear in V_G rather than the common range used in Fig. 4. The current response in this regime clearly shows the stability of both on and off states, and enables accurate determination of the protonic conductivity (see Supplementary information). Reduced fidelity at $f > 10$ Hz is due limited proton conductivity and the $\sim 4 \mu\text{m}$ length of the PEO gate dielectric strip.

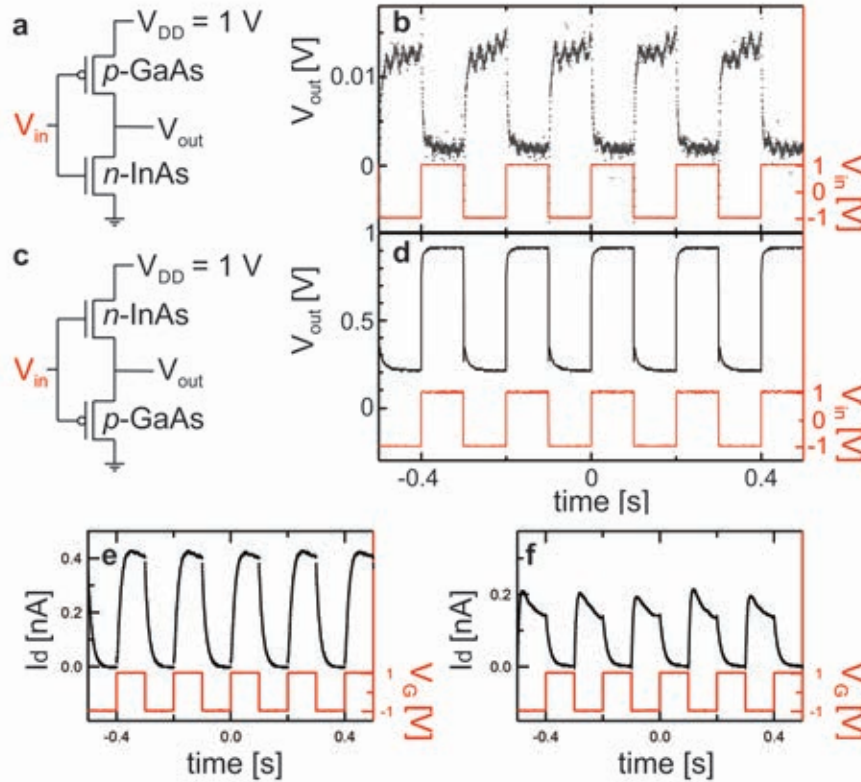


Figure 4: **Demonstration of hybrid organic/inorganic complementary circuitry for logic and signal amplification** **a** Schematic of a traditional inverter circuit featuring undoped PEO nanowire transistors with p -GaAs and n -InAs channels. **b** Plot of input voltage V_{in} (red, right axis) and output voltage V_{out} (black, left axis) vs time for the inverter circuit in **a**. The operating frequency $f = 5$ Hz and the supply voltage $V_{DD} = +1$ V. This circuit embodies the logical NOT operation. **c** similar circuit to **a** but with the p -GaAs and n -InAs transistors swapped to demonstrate the effect of high contact resistance in the p -GaAs device. **d** Plot of input voltage V_{in} (red, right axis) and output voltage V_{out} (black, left axis) vs time for the complementary circuit in **c**. The operating frequency $f = 5$ Hz and the supply voltage $V_{DD} = +1$ V. Device optimisation could see this circuit used to achieve signal amplification. **e/f** Plots of channel current I_d (black, left axis) and gate voltage V_G (red, right axis) for the individual **e** n -InAs and **f** p -GaAs transistors used in **a/c**. Both are presented over the common $-1 < V_G < +1$ V operating range used in complementary circuit operation, as opposed to the linear regime used in Fig. 3.

Correspondingly the on-state stability improves at higher f due to the reduced time available between V_G switching events. The other notable feature in Fig. 3 is the gradual loss in square-wave fidelity with increased frequency, due to the PEO’s limited proton conductivity. The loss of flat maxima/minima for $f \geq 10$ Hz indicates that the time between pulses has dropped below the time required for complete charge separation across the PEO strip. This separation time depends on the PEO strip length but also the proton concentration n and proton mobility μ (see Supplementary Information). Thus, our choice of long PEO strips here is deliberate – it enables us to accurately measure proton conductivity $\sigma = ne\mu$, where e is the proton charge. These strips could be shortened in future devices. We return to proton conductivity in more detail in the Discussion section below.

Hybrid protonic-electronic complementary circuit architectures. A major goal in bioelectronics is to build amplifiers and logic circuits featuring proton-to-electron transducers. We can achieve this using complementary circuit architectures given our approach provides both n - and p -channel transistors. Figure 4 demonstrates two functional circuits featuring p -GaAs and n -InAs transistors with undoped PEO gate dielectrics. We mimicked a wet biological environment by opening the H_2O reservoir to the chamber for 5 – 10 mins – much longer than the 10 – 20 s required in previous experiments – to facilitate H_2O condensation on the sample surface and complete hydration of the PEO films. Cross-linking of the PEO to the substrate during the electron-beam patterning process ensures the nanoscale strips remain stable when immersed in H_2O and PBS solutions for periods exceeding many hours.^{19,30} The first circuit is the traditional inverter circuit in Fig. 4a. The transistor channels were connected in series with the p -GaAs drain at $V_{DD} = 1$ V and the n -InAs source at ground. Here it is essential that the n -type and p -type transistors have strongly overlapping operating gate voltage ranges because the gate electrodes of both transistors must be fed a common electronic square-wave V_{in} (Fig. 4b - red) for complementary operation. The individual n - and p -type transistor responses over the common input voltage range are presented in Figs. 4e/f. The electronic output voltage V_{out} in Fig. 4b (black) was extracted between the

transistors and found to be inverted relative to V_{in} , effectively embodying the logical NOT operation. The limited V_{out} in Fig. 4b is due to the high p -GaAs contact resistance; V_{out} ideally swings between V_{DD} and ground. We demonstrate this using the circuit in Fig. 4c where the p -GaAs and n -InAs devices are reversed. Figure 4d shows V_{out} (black trace) now reaching close to V_{DD} when $V_{in} = 1\text{V}$ (red trace). The high p -GaAs contact resistance prevents V_{out} reaching ground. The signal fidelity is much better and although the gain is less than one, with some optimisation it should be possible to achieve amplification with this circuit (this configuration is commonly known as a complementary push-pull amplifier). We note here the integral role played by the nanoscale patterned, proton-conducting PEO in enabling detection circuits capable of performing logic and amplification. The key is that the PEO can ultimately perform the initial signal detection with nanoscale precision. Thereafter the signal is distributed via proton transport to arrive simultaneously at the n - and p -nanowires. The simultaneous signal arrival is vital to the behavior in Fig. 4; high-fidelity amplification and/or logic cannot be obtained if the signals arrive at different times. Using a biocompatible protonic conductor such as PEO with a well-defined path length as the link between the detection site and the two transistors ensures the circuit functions correctly. Additional steps and challenges towards monolithic integration of our transducing circuits are addressed in the Supplementary Information.

Discussion

Protonic conductivity of nanoscale strips of undoped PEO. Improving signal fidelity and switching speed can be done by optimising device geometry, increasing proton conductivity, or both. As an important step towards this, we measured the protonic conductivity by utilising the I_d transient response of single devices since it is governed by proton drift and follows the exponential form expected for an RC circuit.²⁰ Using our knowledge of the PEO strip dimensions and capacitance, we obtain a proton conductivity $\sigma \sim 10^{-7}$ S/cm (see

Supplementary Information for full details, fitting data and σ versus P_{rel}). Interestingly, the proton conductivity for our undoped PEO films is two to three orders of magnitude higher than previously reported for undoped PEO thin films measured in a parallel-plate electrode geometry.^{21,31} We attribute this to two key differences between our samples and prior studies.

Firstly, the nanoscale patterned PEO in our device geometry is highly exposed, has a tiny volume and a very high surface-to-volume ratio. By contrast, the parallel-plate geometry used in previous works is well known to strongly inhibit water uptake because water can only access via the small exposed side-areas of the thin-film. This can lead to significant inaccuracies (under-measurement) in hydration-dependent studies.^{23,32} As a result, we expect much higher real H^+/OH^- concentration for our devices compared to previous works – and a corresponding increase in σ – under comparable atmospheric conditions. Note that σ may also be influenced by increased μ arising from the swelling behaviour of PEO upon water uptake,¹⁹ however we believe our data is best explained by increasing n with increased P_{rel} as detailed in the Supplementary Information. The enhanced efficacy of water uptake for our nanoscale PEO strips is an advantage for our devices because improved protonic conductivity brings improved frequency response, irrespective of whether higher σ comes about from increased n or μ . The more traditional capacitor geometry is favoured instead for relative ease of fabrication and the precision with which the geometric parameters can be obtained. Note however that geometric parameter uncertainty for our device is insufficient to account for the much higher conductivity we observe; we can measure the dimensions of our PEO strips accurately using atomic force microscopy.

Secondly, in prior work care was taken to exclude water during fabrication for accurate dry-state measurements by, e.g., purifying the as-bought chemicals, carefully dehydrating them and conducting processing in an N_2 environment. We took no such steps to remove residual H_2O or exclude H_2O from the processing, indeed, it is used as the developer solution in our PEO electron-beam lithography process.¹⁰ Taken together, the device geometry and increased exposure to H_2O – whether at development or simply by exposing the films to

ambient – explains well the comparatively high H^+/OH^- concentration in our PEO films. PEO has a known strong affinity for adsorbed H_2O , which is difficult to remove without taking the temperature well above the PEO melting point under high vacuum conditions.²¹ We cannot use this approach as it destroys our nanopatterned strips. We find strong proton conduction is retained even if our sample is held at high vacuum for over 15 hours at room temperature, consistent with earlier work by Binks & Sharples.²¹ This is advantageous from an applied perspective as it improves stability against environmental changes during device operation. The ability to exclude salt-doping without sustaining performance losses is highly advantageous in terms of biocompatibility and fabrication simplicity. A potential issue with the use of InAs and GaAs nanowires is toxicity due to leaching of their surface oxides, which can be unstable under biological conditions. This can be mitigated by simple surface treatments to stabilize the oxide^{33–35} but coverage by PEO may also aid with this. Additionally, signal transduction via the nano-patterned PEO strip means we can physically separate the biological detection site from the nanowires, enabling encapsulation of the nanowire/PEO interface with, e.g., parylene, to ensure fully biocompatible circuit design.

We briefly comment on why we believe ionogenic impurities, e.g., trace salt, play no significant role in our undoped PEO. Firstly, supplier assays for our PEO and solvents have salt contamination at $< 0.3\%$, which we have confirmed via our own inductively coupled plasma mass spectrometry analyses. Secondly, we can support this by noting that equivalent devices featuring PEO films doped at 9–12.5% level with LiClO_4 , as is conventional for PEO-based polymer dielectrics,^{12,20,31,36,37} give approximately equivalent electrical performance (see Supplementary Figure S2/3 and associated discussion). For sake of argument, let us momentarily assume contamination by a non-Li salt, e.g., Na^+ or K^+ . These ions *a priori* should have similar mobility as they are subject to the same chain relaxation transport mechanism³⁶ as H^+ or Li^+ . In reality, the mobility would be slightly lower by virtue of the higher ionic mass/size. For our undoped PEO to give ionic conductivity/performance comparable to Li-doped PEO, without protonic transport playing any role, we would require

a similar ion concentration, and thus salt contamination at $\sim 10\%$ level – clearly this is inconsistent with our chemical analysis. The earlier work by Binks & Sharples arrived at an identical conclusion, namely that protonic transport and not ionogenic impurity effects explain the conductivity behaviour of undoped PEO with significant water content.²¹

The path to improved frequency performance. The data presented here highlight the potential for nanowires in amplifying and processing biological signals; they also suggest points for enhancing our design. A key limitation presently is the frequency response. Our knowledge of the proton conductivity allows us to quantify potential improvements on this aspect. Shortening the PEO strip would be the simplest action but increasing its width to cover as much of the nanowire as possible without touching the contacts would also help. If we assume the PEO strip length can be reduced to $l \sim 100$ nm with $\sigma = 10^{-7}$ S/cm, we would expect switching speeds exceeding 1.6 kHz, more than sufficient for neural interfacing applications where signal frequencies are typically below 200 Hz.³ Note also that our devices inevitably involve a two-stage electron-to-proton-to-electron transduction process since an electrical voltage signal V_G applied to the gate electrode drives H^+/OH^- transport which in turn affects I_d in the nanowire. Faster performance may be expected if using PEO-coated nanowires as a direct protonic sensing element. That is, V_{in} could arise from direct protonic/ionic coupling of the PEO to a biological system of interest, e.g., a neuron,³ biological nanopore,³⁸ or protonic transistor channel,^{4,5} rather than the gate electrode used for our basic demonstration here. This will be a subject of future work. Further improvements in conductivity and frequency response may also be possible by doping the PEO with additional H^+ ,³⁹ or by substituting the PEO for a material with higher protonic conductivity such as maleic chitosan.^{4,9} A critical requirement and challenge for any such substitute material is demonstrating nanoscale patternability. As discussed above, the patternability of PEO is an essential feature of our design as it enables simultaneous sensing of the biological by the n - and p -type nanowire channels.

Conclusion. In conclusion, proton conduction in nanopatterned PEO films can be used

for highly-effective gating of both *n*-type InAs and *p*-type GaAs nanowire transistors. The proton conductivity for our undoped PEO is 2–3 orders higher than previous reports,³¹ likely from differences in device architecture and processing. We demonstrated proton-to-electron signal transduction with good fidelity at frequencies up to 50 Hz. Gains in operating speed can in future be achieved by reducing the PEO thickness or doping the PEO with excess H⁺.^{24,40} We also demonstrated basic hybrid organic/inorganic complementary circuit architectures including a prototype inverter and repeater/push-pull amplifier circuit, although gain will require lower contact resistance to our *p*-GaAs nanowires. Our work provides a clear path towards high fidelity proton-to-electron transduction devices and circuits enabling functions, e.g., logic and amplification, using nanowire transistors featuring nanoscale proton conducting elements.

Methods

Device Fabrication. Nanowires with lengths 3 – 6 μm and diameters 50 nm or 100 nm were grown by either vapour-liquid-solid MOCVD or MBE for InAs and GaAs, respectively. Surface accumulation of electrons in nominally undoped InAs nanowires results in *n*-type conduction without doping. The *p*-GaAs nanowires were Be-doped with acceptor concentration approximately $N_A = 1 \times 10^{18} \text{ cm}^{-3}$. The N_A value is expected to closely match that for planar growth, although some variation may occur along the radial axis.⁴¹ After transferring the nanowires to the measurement substrate, source (S) and drain (D) contacts and gate electrodes (G1 and G2) were defined by electron beam lithography using a polymethylmethacrylate (PMMA) resist. The contacts and electrodes consisted of 25/75 nm Ni/Au for InAs and 200 nm AuBe for GaAs deposited by thermal evaporation. For InAs, passivation in an aqueous $(\text{NH}_4)_2\text{S}_x$ solution immediately prior to metal deposition ensured ohmic contacts. For GaAs, the native oxide was removed by a 30 s HCl etch prior to metal deposition and post-deposition annealing at 300°C for 30 s caused the AuBe to diffuse into

the GaAs and form ohmic contacts. Polymer electrolyte thin films were formed by dissolving dry PEO, spin coating the solution onto the sample, and evaporating the methanol by baking at 90°C for 30 mins. This is similar to the process described in Refs. 12,37. The thin films were patterned directly by electron beam lithography in a Raith 150-Two system using an accelerating voltage of 5 kV, dose of 100 – 150 $\mu\text{C}/\text{cm}^2$ and typical beam current ~ 15 pA. Extended details on electron beam patterning PEO appear in Ref. 10.

Electrical Characterization. Electrical measurements were performed under a hydration controlled atmosphere in a specially constructed vacuum chamber. The chamber was connected to a reservoir of highly purified water isolated by a bleed valve. The water was deionised Milli-Q water (18 M Ω) that had been degassed with three freeze-pump-thaw cycles. Opening the bleed valve admitted water vapour to the chamber, increasing chamber pressure P in discrete steps from vacuum ($P = 0$ mbar, $P_{rel} = 0$) to saturated vapour pressure ($P = 24$ mbar, $P_{rel} = 1$) as described previously^{22,32} and in detail in the Supplementary Information. A Keithley K2450 was used to supply $V_{sd} = 2$ mV (n -InAs) or $V_{sd} = 1$ V (p -GaAs) and measure I_d . To obtain the quasi-dc characteristics in Fig. 2 an additional Keithley K2400 was used to supply V_G and monitor the gate leakage current, which remained below the noise floor (< 0.1 nA) throughout the experiments. V_G was incremented at 5 mV/s. For ac measurements (Figs. 3 and 4) a H.P. 33120A function generator supplied the square wave voltage to the gate while I_d was amplified by a Femto DLPCA-200 low noise current amplifier and recorded using an Agilent DSO-X 3024A digital oscilloscope. For the complementary circuits in Fig. 4b/d, V_{out} was measured directly by the oscilloscope without external amplification. An extended discussion of these circuits and the methods used to measure proton conductivity is given in the Supplementary Information.

Supporting Information. Full details on device fabrication, comparison of devices with LiClO₄-doped and undoped PEO, full humidity dependence data, full protonic conductivity data, and comments on transient response and issues related to full monolithic integration of complementary circuits are included in Supporting Information. This material is available

free of charge via the Internet at <http://pubs.acs.org>.

Author Contributions. DJC, ABM, PM and APM conceived and directed the project. DJC, RU and AMB fabricated devices. DJC and ABM performed measurements and analysed data. HJJ, HHT and CJ provided InAs nanowires. PK and JN provided GaAs nanowires. DJC and APM wrote the manuscript, assisted by ABM and PM. All authors commented on the manuscript.

Corresponding author. *E-mail: adam.micolich@nanoelectronics.physics.unsw.edu.au

Acknowledgement

This work was funded by the Australian Research Council (ARC), the University of New South Wales, the University of Queensland, Danish National Research Foundation and the Innovation Fund. APM acknowledges an ARC Future Fellowship (FT0990285) and DJC acknowledges Australian Nanotechnology Network Short Term Visit support. PM is an ARC Discovery Outstanding Research Award Fellow and the work at UQ was funded under the ARC Discovery Program (DP140103653). The Centre for Organic Photonics and Electronics is a strategic initiative of the University of Queensland. We thank Helen Rutledge for conducting the inductively coupled plasma mass spectrometry measurements. This work was performed in part using the NSW and ACT nodes of the Australian National Fabrication Facility (ANFF) and the Mark Wainwright Analytical Centre at UNSW.

References

- (1) Berggren, M.; Richter-Dahlfors, A. *Advanced Materials* **2007**, *19*, 3201–3213.
- (2) Owens, R. M.; Malliaras, G. G. *MRS bulletin* **2010**, *35*, 449456.
- (3) Khodagholy, D.; Doublet, T.; Quilichini, P.; Gurfinkel, M.; Leleux, P.; Ghestem, A.; Ismailova, E.; Hervé, T.; Sanaur, S.; Bernard, C.; Malliaras, G. G. *Nature Communications* **2013**, *4*, 1575.

- (4) Zhong, C.; Deng, Y.; Roudsari, A. F.; Kapetanovic, A.; Anantram, M.; Rolandi, M. *Nature Communications* **2011**, *2*, 476.
- (5) Ordinario, D. C.; Phan, L.; Walkup IV, W. G.; Jocson, J. M.; Emil, K.; Nina, H.; A., G. A. *Nature Chemistry* **2014**, *6*, 596602.
- (6) Strakosas, X.; Bongo, M.; Owens, R. M. *Journal of Applied Polymer Science* **2015**, *132*, 41735.
- (7) Tybrandt, K.; Larsson, K. C.; Agneta, R.-D.; Berggren, M. *Proceedings of the National Academy of Sciences U.S.A.* **2012**, *107*, 9929–9932.
- (8) Tybrandt, K.; Forchheimer, R.; Berggren, M. *Nature Communications* **2012**, *3*, 871.
- (9) Deng, Y.; Josberger, E.; Jin, J.; Rousdari, A. F.; Helms, B. A.; Zhong, C.; Anantram, M. P.; Rolandi, M. *Scientific Reports* **2013**, *3*, 2481.
- (10) Carrad, D. J.; Burke, A. M.; Lyttleton, R. W.; Joyce, H. J.; Tan, H. H.; Jagadish, C.; Storm, K.; Linke, H.; Samuelson, L.; Micolich, A. P. *Nano Letters* **2014**, *14*, 94–100.
- (11) Cho, J. H.; Lee, J.; Xia, Y.; Kim, B.; He, Y.; Renn, M. J.; Lodge, T. P.; Daniel Frisbie, C. *Nature Materials* **2008**, *7*, 900–906.
- (12) Liang, D.; Gao, X. P. A. *Nano Letters* **2012**, *12*, 3263–3267.
- (13) Cui, Y.; Wei, Q.; Park, H.; Lieber, C. M. *Science* **2001**, *293*, 1289–1292.
- (14) Misra, N.; Martinez, J. A.; Huang, S. J.; Wang, Y.; Stroeve, P.; Grigoropoulos, C. P.; Noy, A. *Proceedings of the National Academy of Sciences* **2009**, *106*, 1378013784.
- (15) Du, J.; Liang, D.; Tang, H.; Gao, X. P. *Nano Letters* **2009**, *9*, 4348–4351.
- (16) Upadhyay, S.; Frederiksen, R.; Lloret, N.; De Vico, L.; Krogstrup, P.; Jensen, J. H.; L., M. K.; J., N. *Applied Physics Letters* **2014**, *104*, 203504.

- (17) Tomioka, K.; Yoshimura, M.; Fukui, T. *Nature* **2012**, *488*, 189–192.
- (18) Schmid, H.; Borg, M.; Moselund, K.; Gignac, L.; Breslin, C. M.; Bruley, J.; Cutaia, D.; Riel, H. *Applied Physics Letters* **2015**, *106*, 233101.
- (19) Krsko, P.; Sukhishvili, S.; Mansfield, M.; Clancy, R.; Libera, M. *Langmuir* **2003**, *19*, 5618–5625.
- (20) Kim, S. H.; Hong, K.; Xie, W.; Lee, K. H.; Zhang, S.; Lodge, T. P.; Frisbie, C. D. *Advanced Materials* **2013**, *25*, 1822–1846.
- (21) Binks, A. E.; Sharples, A. *Journal of Polymer Science Part A-2: Polymer Physics* **1968**, *6*, 407420.
- (22) Mostert, A. B.; Davy, K. J. P.; Ruggles, J. L.; Powell, B. J.; Gentle, I. R.; Meredith, P. *Langmuir* **2010**, *26*, 412–416.
- (23) Mostert, A. B.; Powell, B. J.; Gentle, I. R.; Meredith, P. *Applied Physics Letters* **2012**, *100*, 093701–093701–3.
- (24) DeLongchamp, D. M.; Hammond, P. T. *Langmuir* **2004**, *20*, 5403–5411.
- (25) Roddaro, S.; Nilsson, K.; Astromskas, G.; Samuelson, L.; Wernersson, L.; Karlström, O.; Wacker, A. *Applied Physics Letters* **2008**, *92*, 253509.
- (26) Burke, A. M.; Waddington, D. E. J.; Carrad, D. J.; Lyttleton, R. W.; Tan, H. H.; Reece, P. J.; Klochan, O.; Hamilton, A. R.; Rai, A.; Reuter, D.; Wieck, A. D.; Micolich, A. P. *Physical Review B* **2012**, *86*, 165309.
- (27) Carrad, D. J.; Burke, A. M.; Reece, P. J.; Lyttleton, R. W.; Waddington, D. E. J.; Rai, A.; Reuter, D.; Wieck, A. D.; Micolich, A. P. *Journal of Physics: Condensed Matter* **2013**, *25*, 325304.

- (28) Egard, M.; Johansson, S.; Johansson, A.; Persson, K.; Dey, A. W.; Borg, B. M.; Thelander, C.; Wernersson, L.; Lind, E. *Nano Letters* **2010**, *10*, 809–812.
- (29) Lind, E.; Persson, A. I.; Samuelson, L.; Wernersson, L. *Nano Letters* **2006**, *6*, 1842–1846.
- (30) Hong, Y.; Krsko, P.; Libera, M. *Langmuir* **2004**, *20*, 11123–11126.
- (31) Fullerton-Shirey, S. K.; Maranas, J. K. *Macromolecules* **2009**, *42*, 2142–2156.
- (32) Mostert, A. B.; Powell, B. J.; Pratt, F. L.; Hanson, G. R.; Sarna, T.; Gentle, I. R.; Meredith, P. *Proceedings of the National Academy of Sciences* **2012**, *109*, 8943–8947.
- (33) Cho, Y.; Ivanisevic, A. *NanoBiotechnology* **2006**, *2*, 51–59.
- (34) Jewett, S. A.; Ivanisevic, A. *Accounts of Chemical Research* **2012**, *45*, 1451–1459.
- (35) Park, G. et al. *Advanced Healthcare Materials* **2014**, *3*, 515–525.
- (36) Gorecki, W.; Belorizky, E.; Berthier, C.; Donoso, P.; Armand, M. *Electrochimica Acta* **1992**, *37*, 1685–1687.
- (37) Panzer, M. J.; Newman, C. R.; Frisbie, C. D. *Applied Physics Letters* **2005**, *86*, 103503.
- (38) Noy, A. *Advanced Materials* **2011**, *23*, 807–820.
- (39) Hashmi, S. A.; Kumar, A.; Maurya, K. K.; Chandra, S. *Journal of Physics D: Applied Physics* **1990**, *23*, 1307.
- (40) Ratner, M. A.; Shriver, D. F. *Chemical Reviews* **1988**, *88*, 109–124.
- (41) Casadei, A.; Krogstrup, P.; Heiss, M.; Röhr, J. A.; Colombo, C.; Ruelle, T.; Upadhyay, S.; Sørensen, C. B.; Nygård, J.; Fontcuberta i Morral, A. *Applied Physics Letters* **2013**, *102*, 013117.