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Complex geometry of quantum cones

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The algebras obtained as fixed points of the action of the cyclic group \mathbb{Z}_N on the coordinate algebra of the quantum disc are studied. These can be understood as coordinate algebras of quantum or non-commutative cones. The following observations are made. First, contrary to the classical situation, the actions of \mathbb{Z}_N are free and the resulting algebras are homologically smooth. Second, the quantum cone algebras admit differential calculi that have all the characteristics of calculi on *smooth* complex curves. Third, the corresponding volume forms are exact, indicating that the constructed algebras describe manifolds with boundaries.

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1 Introduction

Recent studies of deformed orbifolds or manifolds with singularities revealed an intriguing phenomenon: many classical surfaces Σ with singular points admit non-commutative or q-versions Σ_q that are smooth. Smoothness can be understood on several levels. Most naively, one observes that the multiple roots of the defining polynomial relations become separated upon the quantization [6], [3]. Furthermore, the classically non-free action of a finite group on a smooth surface that gives rise to the singularity becomes free once the surface and the action are quantized [4]. On a more sophisticated level, the deformed coordinate algebras $\mathcal{O}(\Sigma_q)$ are homologically smooth, i.e. they fit into a finite exact sequence of finitely generated and projective $\mathcal{O}(\Sigma_q)$ -bimodules [4]. Often, $\mathcal{O}(\Sigma_q)$ is a (twisted) Calabi-Yau algebra of dimension two, i.e $\mathcal{O}(\Sigma_q)$ is not only homologically smooth, but also its Hochschild cohomology with values in the enveloping algebra of $\mathcal{O}(\Sigma_q)$ is trivial except in degree two, where it is isomorphic to $\mathcal{O}(\Sigma_q)$ with the bimodule structure twisted by an automorphism [15]. This last observation can be interpreted as an existence of the volume form, and is tantamount to the Poincaré duality [21].

Many two-dimensional orbifolds are obtained by glueing cones, hence quantum cones seem to be a natural point to start. In this note we carry out parts of the programme of differential smoothing initiated in [9] and discuss differential structures on quantum cones that allow one to view them as complex curves. We define coordinate algebras of quantum cones $\mathcal{O}(C_{q,\gamma}^N)$ as fixed points of the action of cyclic groups \mathbb{Z}_N on the non-commutative disc algebra $\mathcal{O}(D_{q,\gamma})$ [13]. This action turns out to be free and $\mathcal{O}(C_{q,\gamma}^N)$ are (twisted) Calabi-Yau algebras. We derive differential calculi on $\mathcal{O}(C_{q,\gamma}^N)$ from that on $\mathcal{O}(D_{q,\gamma})$ and show that they define holomorphic structures on $\mathcal{O}(C_{q,\gamma}^N)$ in the sense of [12]. The resulting volume forms on $\mathcal{O}(C_{q,\gamma}^N)$ are exact, which might indicate that quantum cones have boundaries.

2 Quantum cones are smooth

The coordinate algebra $\mathcal{O}(D_{q,\gamma})$ of the quantum disc is defined as the complex *-algebra generated by z,z^* subject to relation

$$z^*z - qzz^* = \gamma, (1)$$

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where $q, \gamma \in \mathbb{R}$, q > 0 are parameters. This is obviously a deformation of the plane (or the unit disc), which corresponds to q = 1, $\gamma = 0$. $\mathcal{O}(D_{q,0})$ is the quantum or Manin's plane, and the case $\mathcal{O}(D_{1,1})$ is the quantum oscillator algebra. In fact, if $\gamma > 0$, then the generators can be rescaled, so that γ can be eliminated from (1). If $q \in (0,1)$, the algebra $\mathcal{O}(D_{q,1-q})$ can be completed to the C^* -algebra isomorphic to the Toeplitz algebra [13]. For an accessible introduction to geometry of the quantum disc we refer to [20, Chapter 1].

 $\mathcal{O}(D_{q,\gamma})$ admits an action of the cyclic group \mathbb{Z}_N , which can be described in terms of the \mathbb{Z}_N -grading on generators as $\deg(z)=1$, $\deg(z^*)=N-1=-1$ mod N. The fixed point subalgebra of this action or the degree-zero part of $\mathcal{O}(D_{q,\gamma})$ is generated by $a=zz^*$, $b=z^N$, $b^*=z^{*N}$. These satisfy the relations

$$ab = q^{N}ba + \gamma[N]_{q}b, \qquad bb^{*} = \prod_{l=0}^{N-1} (q^{-l}a + \gamma[-l]_{q}), \qquad b^{*}b = \prod_{l=1}^{N} (q^{l}a + \gamma[l]_{q}),$$
 (2)

where $[n]_q=(1-q^n)/(1-q)$ are q-integers. The complex *-algebra generated by a selfadjoint a and b, b^* subject to relations (2) is called the *coordinate algebra of the quantum N-cone* and is denoted by $\mathcal{O}(C_{q,\gamma}^N)$. In case $q=1,\ \gamma=0,\ a,\ b,\ b^*$ can be thought of as coordinates and then relations (2) describe a cone in $\mathbb{C}\times\mathbb{R}$, hence the name. The polynomials in a on the right-hand sides of the last two equations in (2) have no repeated roots as long as $\gamma\neq 0$. One might therefore expect that in this case the algebras $\mathcal{O}(C_{q,\gamma}^N)$ describe smooth non-commutative surfaces.

An action of a (finite Abelian) group G on a quantum space Σ_q is free if and only if the corresponding coordinate algebra $\mathcal{O}(\Sigma_q)$ is $strongly\ G$ -graded. This means that for all $g\in G$, one can find finite number of elements X_i of G-degree g and Y_i of degree g^{-1} such that $\sum_i X_i Y_i = 1$. In the case of the cyclic group \mathbb{Z}_N suffices it to find such X_i of degree 1 and 10 degree 11, see [17, Section AI.3.2].

Theorem 2.1 For all $\gamma \neq 0$, the \mathbb{Z}_N -action on $D_{q,\gamma}$ is free. Consequently, $\mathcal{O}(C_{q,\gamma}^N)$ are homologically smooth algebras for all values of N > 1, q and $\gamma \neq 0$.

Proof. Set $X_0 = \alpha_0 z^{*N-1}$, $X_i = \alpha_i a^{i-1} z$, $Y_0 = z^{N-1}$ and $Y_i = z^*$, $i = 1, \dots, N-1$, where all the $\alpha_r \in \mathbb{C}$ are to be determined. By equations (1),

$$z^{*N-1}z^{N-1} = \prod_{l=1}^{N-1} (q^l a + \gamma[l]_q) =: \sum_{r=0}^{N-1} \beta_r a^r.$$

The comparison of the coefficients of the powers of a in the condition $\sum_{i=0}^{N-1} X_i Y_i = 1$, yields the system of N equations for the α_r : $\alpha_0 \beta_0 = 1$ and $\alpha_0 \beta_i + \alpha_i = 0$, for $i = 1, 2, \dots, N-1$, whose determinant is $\beta_0 = \gamma^{N-1} \prod_{l=1}^{N-1} [l]_q \neq 0$. Hence the required X_i and Y_i can be found.

Since the enveloping algebra of $\mathcal{O}(D_{q,\gamma})$ is left Noetherian of finite global dimension and since the \mathbb{Z}_N -actions defining the cones are free if $\gamma \neq 0$, the algebras $\mathcal{O}(C_{q,\gamma}^N)$ are homologically smooth by [14, Corollary 6] or [4, Criterion 1].

The freeness of the action of \mathbb{Z}_N on $D_{q,\gamma}$ is equivalent to the statement that the quantum disc is the total space of a \mathbb{Z}_N -principal quantum bundle over the quantum cone $C_{q,\gamma}^N$ [8], [7]. This bundle is non-trivial, i.e. $\mathcal{O}(D_{q,\gamma})$ is not isomorphic to $\mathcal{O}(C_{q,\gamma}^N) \otimes \mathbb{C}\mathbb{Z}_N$ as a \mathbb{Z}_N -graded left $\mathcal{O}(C_{q,\gamma}^N)$ -module. The triviality would require the existence of invertible elements in $\mathcal{O}(D_{q,\gamma})$ of degree n, for all $n \in \mathbb{Z}_N$. However, only non-zero scalar multiples of the identity are invertible in $\mathcal{O}(D_{q,\gamma})$, and these have degree zero.

Theorem 2.2 For all $\gamma \neq 0$, $q \neq 1$, $\mathcal{O}(C_{q,\gamma}^N)$ are (twisted) Calabi-Yau algebras.

Proof. Define $x := 1 + \gamma^{-1}(q-1)a$. The relations (2) yield

$$xb = q^N bx, bb^* = \left(\frac{\gamma}{1-q}\right)^N \prod_{l=0}^{N-1} (1-q^{-l}x), b^*b = \left(\frac{\gamma}{1-q}\right)^N \prod_{l=1}^N (1-q^l x).$$
 (3)

Since the change of variables $a\mapsto x$ is reversible, equations (3) simply give a different presentation of $\mathcal{O}(C^N_{q,\gamma})$. This presentation allows one to view $\mathcal{O}(C^N_{q,\gamma})$ as a *generalized Weyl algebra* over the polynomial ring in the variable x [1]. Since $q\neq 1$, the polynomials on the right hand side of (3) have no repeated roots and $\mathcal{O}(C^N_{q,\gamma})$ are twisted Calabi-Yau algebras by [15, Theorem 4.5].

An inspection of relations (3) makes it clear that, for $\gamma > 0$, $q \in (0,1)$ and N odd, the algebras $\mathcal{O}(C_{q,\gamma}^N)$ are isomorphic to coordinate algebras of quantum real weighted projective spaces $\mathcal{O}(\mathbb{RP}_q^2(N;+))$ [3].

3 Differential and complex geometry of quantum cones

By a differential calculus over an algebra A we mean a pair $(\Omega(A),d)$, where $\Omega(A)=\oplus_{n\in\mathbb{N}}\Omega^n(A)$ is an \mathbb{N} -graded algebra such that $\Omega^0(A)=A$, and $d:\Omega^n(A)\to\Omega^{n+1}(A)$ is the degree-one linear map that satisfies the graded Leibniz rule, is such that $d\circ d=0$ and, for all $n\in\mathbb{N}$, $\Omega^n(A)=Ad(A)d(A)\cdots d(A)$ (d(A)) appears n-times). If A is a complex *-algebra, then it is often requested that $\Omega(A)$ be a *-algebra and that 00 and 01. In this situation one refers to 01 as to a *-differential calculus. In this case, following [12], a complex structure on 01 is the bi-grading decomposition of 01,

$$\Omega^{n}(A) = \bigoplus_{k+l=n} \Omega^{(k,l)}(A),$$

such that $*: \Omega^{(k,l)}(A) \to \Omega^{(l,k)}(A)$, and the decomposition $d = \delta + \bar{\delta}$ into differentials $\delta: \Omega^{(k,l)}(A) \to \Omega^{(k+1,l)}(A)$, $\bar{\delta}: \Omega^{(k,l)}(A) \to \Omega^{(k,l+1)}(A)$, such that $\delta(a)^* = \bar{\delta}(a^*)$, for all $a \in \Omega(A)$.

To construct a complex structure on $\mathcal{O}(C_{q,\gamma}^N)$ start with the differential *-calculus $\Omega(D_{q,\gamma})$ on $\mathcal{O}(D_{q,\gamma})$ generated by one-forms dz and dz^* subject to relations

$$zdz = q^{-1}dzz, z^*dz = qdzz^*, dz \wedge dz^* = -q^{-1}dz^* \wedge dz, dz \wedge dz = 0; (4)$$

see e.g. [19]. Note that $\Omega(D_{q,\gamma})=\mathcal{O}(D_{q,\gamma})\oplus\Omega^1(D_{q,\gamma})\oplus\Omega^2(D_{q,\gamma})$, where $\Omega^1(D_{q,\gamma})$ is a free left and right $\mathcal{O}(D_{q,\gamma})$ -module with free generators dz and dz^* , and $\Omega^2(D_{q,\gamma})$ is a free $\mathcal{O}(D_{q,\gamma})$ -module generated by $dz\wedge dz^*$. We define the differential *-calculus $\Omega(C_{q,\gamma}^N)$ on $\mathcal{O}(C_{q,\gamma}^N)$ as the differential *-subalgebra of $\Omega(D_{q,\gamma})$ generated by $a=zz^*$, $b=z^N$ and their differentials da, db. Set

$$\omega_0 := dzz^*, \qquad \omega_1 := db = [N]_q z^{N-1} dz, \qquad \omega := dz \wedge dz^*.$$

Theorem 3.1 Assume that $\gamma \neq 0$ and $q \neq 1$. Let $\Omega^{(1,0)}(C_{q,\gamma}^N)$ be the submodule of $\Omega^1(C_{q,\gamma}^N)$ generated by ω_0 , ω_1 , let $\Omega^{(0,1)}(C_{q,\gamma}^N)$ be the submodule of $\Omega^1(C_{q,\gamma}^N)$ generated by ω_0^* , ω_1^* , and let $\Omega^{(1,1)}(C_{q,\gamma}^N) = \Omega^2(C_{q,\gamma}^N)$. Set $\Omega^{(k,l)}(C_{q,\gamma}^N) = 0$ if k > 1 or l > 1. Let δ be the projection of d onto $\Omega^{(1,0)}(C_{q,\gamma}^N)$ and let $\bar{\delta}$ be the projection of d onto $\Omega^{(0,1)}(C_{q,\gamma}^N)$. Then the above bi-grading decomposition of $\Omega(C_{q,\gamma}^N)$ defines a complex structure on $\mathcal{O}(C_{q,\gamma}^N)$. Furthermore, $\Omega^2(C_{q,\gamma}^N) = \mathcal{O}(C_{q,\gamma}^N) = \omega \mathcal{O}(C_{q,\gamma}^N)$.

Proof. The only non-trivial statement to be proven here is that both ω_0 and ω_0^* belong to $\Omega^1(C_{q,\gamma}^N)$ (there sum is in $\Omega^1(C_{q,\gamma}^N)$, since it is equal to da). If this is established, then the bi-graded decomposition of $\Omega^1(C_{q,\gamma}^N)$ and maps δ , $\bar{\delta}$ are well-defined. Also, since $\omega=-d\omega_0$, the last assertion will follow. We start by proving the following

Lemma 3.2 For all positive $q \neq 1$ and positive $n \in \mathbb{N}$, if an ideal in the polynomial ring $\mathbb{C}[w]$ contains both $X_n(w) = \prod_{l=1}^{n-1} (1+q^{-l}w)$ and $Y_n(w) = \prod_{l=1}^{n-1} (1+q^{l}w)$, then it also contains the identity.

Proof. By induction in n. Clearly, the statement is true for n=1,2. Assume it is true for k, i.e. there exist polynomials f(w) and g(w) such that $f(w)X_k(w)+g(w)Y_k(w)=1$. Let J be an ideal containing $X_{k+1}(w)$ and $Y_{k+1}(w)$. Then

$$(1+q^{-k}w)(1+q^kw) = (f(w)X_k(w)+g(w)Y_k(w))(1+q^{-k}w)(1+q^kw)$$

= $f(w)X_{k+1}(w)(1+q^kw)+g(w)Y_{k+1}(w)(1+q^{-k}w) \in J.$

Similarly, by considering $(f(qw)X_k(qw)+g(qw)Y_k(qw))(1+qw)(1+q^{-k+1}w)(1+q^{-k}w)$ one finds that $(1+qw)(1+q^{-k}w)(1+q^{-k+1}w)\in J$, and hence also $(1+q^{-1}w)(1+q^kw)(1+q^{k-1}w)\in J$, by symmetry. By comparing powers of w, one easily finds that there exist complex numbers $\alpha_0,\alpha_1,\alpha_2$ such that $(\alpha_0+\alpha_1w)(1+q^kw)+\alpha_2(1+qw)(1+q^{-k+1}w)=1$. Multiplying this by $1+q^{-k}w$ we find that $1+q^{-k}w\in J$. Again, by symmetry $1+q^kw\in J$, and clearly there exists a linear combination of these that makes up 1.

Define $w := \gamma^{-1}(1-q)qa - q$. Using relations (1) and (4) one finds

$$b^*db = q[N]_q \left(\frac{\gamma}{1-q}\right)^{N-1} \prod_{l=1}^{N-1} (1+q^l w) \,\omega_0, \quad bdb^* = [N]_q \left(\frac{q^{-1}\gamma}{1-q}\right)^{N-1} \prod_{l=1}^{N-1} (1+q^{-l}w) \,\omega_0. \tag{5}$$

Hence $\prod_{l=1}^{N-1}(1+q^lw)\omega_0, \prod_{l=1}^{N-1}(1+q^{-l}w)\omega_0\in\Omega^1(C_{q,\gamma}^N)$ and, by Lemma 3.2, $\omega_0\in\Omega^1(C_{q,\gamma}^N)$. Since the calculus is closed under the *-operation also $\omega_0^*\in\Omega^1(C_{q,\gamma}^N)$. Now the theorem follows.

The existence of relations such as (5) means that the generators ω_0 and ω_1 of the left (or right) $\mathcal{O}(C_{q,\gamma}^N)$ -module $\Omega^{(1,0)}(C_{q,\gamma}^N)$ are not independent. That $\Omega^{(1,0)}(C_{q,\gamma}^N)$ is a finitely generated and projective $\mathcal{O}(C_{q,\gamma}^N)$ -module may be argued as follows. $\Omega^{(1,0)}(C_{q,\gamma}^N)$ is a submodule of the free $\mathcal{O}(D_{q,\gamma})$ -module $\mathcal{O}(D_{q,\gamma})$ dz, generated by $z^{N-1}dz$ and z^*dz , hence it can be identified with the $\mathcal{O}(C_{q,\gamma}^N)$ -submodule of $\mathcal{O}(D_{q,\gamma})$ generated by z^{N-1} and z^* . This is exactly the $\mathcal{O}(C_{q,\gamma}^N)$ -module consisting of all elements in $\mathcal{O}(D_{q,\gamma})$ of degree N-1. By Theorem 2.1, $\mathcal{O}(D_{q,\gamma})$ is strongly graded, hence any submodule of elements of a fixed degree is projective by [17, A I.3.3 Corollary]. Therefore $\Omega^{(1,0)}(C_{q,\gamma}^N)$ can be seen as module of sections of a non-commutative vector bundle over the quantum cone. Furthermore, since $\mathcal{O}(D_{q,\gamma})$ has no invertible elements of degree 1 or N-1, by the same arguments as those following the proof of Theorem 2.1, the module $\Omega^{(1,0)}(C_{q,\gamma}^N)$ is not free. This is in full agreement with the identification of $\mathcal{O}(C_{q,1-q}^2)$ as the quantum equatorial sphere; holomorphic cotangent bundle over the sphere is not trivial. The employment of the *-conjugation yields similar statements for $\Omega^{(0,1)}(C_{q,\gamma}^N)$.

of the *-conjugation yields similar statements for $\Omega^{(0,1)}(C_{q,\gamma}^N)$.

Perhaps it is no too surprising that the quantum cones $C_{q,\gamma}^N$ admit natural complex structures, for the considerations of Section 2 substantiate the claim that they are smooth quantum manifolds. The exactness of the volume form might indicate that quantum cones correspond to non-commutative manifolds with boundaries. The non-triviality of the cotangent bundle indicates that $C_{q,\gamma}^N$ correspond to non-parallelizable manifolds. For these reasons one might think about the non-commutative smoothing of the cone as of blowing up the singular point into a balloon with a neck rather than as of rounding it.

Finally, one might observe that the differential calculus on the quantum disc studied in here is generated by *q-derivations*. Following this observation, methods and results of [5] can be applied to establish the *differential smoothness* of quantum cones in the sense of [9], i.e. the existence of an isomorphism between complexes of differential and integral forms. The detailed discussion of these issues, however, goes beyond the scope of the present note.

4 Outlook and speculations

My motivation for writing this note was to draw the reader's attention to unexpected results of actions of finite groups on quantum spaces. The classically non-free actions are freed upon quantization, singularities are removed, and the geometric and topological natures of the spaces of fixed points are transformed. This opens up an exciting possibility of deploying the full power of differential geometry in areas in which it could not be utilized previously. Already the study of one of the simplest examples leads to surprises, poses new questions and opens up avenues for further research. An immediate direction of studies would be to try and understand quantum cones as Riemannian surfaces, whether from the spectral point of view [11] or in a more algebraic setup [16], [2]. Once these aspects are understood in sufficient depth, one can attempt

to build physical models based on quantum cones. The lessons thus learnt should be applied to other deformations of orbifolds including those that can be obtained by glueing of quantum cones.

Hitherto we discussed only algebraic aspects of quantum cones, leaving aside topological issues. In view of relations (3), if $\gamma>0$ and $q\in(0,1)$, the C^* -algebra completions of the $\mathcal{O}(C_{q,\gamma}^N)$ to algebras of continuous functions $\mathcal{C}(C_{q,\gamma}^N)$ can be carried out in the same way as for the quantum real weighted projective spaces [3]. Consequently, $\mathcal{C}(C_{q,\gamma}^N)$ is a pullback of N-copies of the Toeplitz algebra along the symbol maps or the algebra of continuous functions on the N-copies of quantum discs glued along their circular boundaries; see [10]. In particular $\mathcal{C}(C_{q,\gamma}^2)$ is isomorphic to the algebra of continuous functions on the quantum equatorial or generic Podleś sphere [18]. This identifies $C_{q,\gamma}^2$ as a deformation of a compact manifold without a boundary. Classically, volume forms for such manifolds are never exact, yet the volume forms described in Theorem 3.1 are exact! The exactness of the volume form is compatible with the interpretation of $C_{q,\gamma}^2$ as the cone but not as the sphere. On the other hand, the interpretation of $C_{q,\gamma}^2$ as a sphere is compatible with the non-triviality of cotangent bundles.

This identification of $C_{q,\gamma}^2$ leads to another question: how is it possible for the sphere to arise from an action of a finite group on the disc? A possible heuristic answer might be to think about the quantum disc as having an internal structure: the quantum disc is composed of layers of quantum discs. The \mathbb{Z}_2 -action does not bulge the disc into a cone, but rather it gently separates two internal layers, which are still joined at the common circle, thus forming a sphere or a balloon. Speculatively, one might interpret the classical circle that the equatorial Podleś sphere contains as the rim of the neck of the balloon (rather than as the equator), and thus reconcile this point of view with the interpretation of the differential smoothing offered at the end of Section 3. For the action of \mathbb{Z}_N , N internal layers are separated.

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